

## INTRODUCTION

**Motivation:** Computing transport coefficients (mobility, thermal conductivity, shear viscosity) in statistical physics

**Goal:** Reduce the variance of the estimator of sensitivity/transport coefficient.

## DYNAMICS

Consider the following family of SDEs with values in  $\mathbb{R}^d$  and additive noise:

$$dX_t^\eta = (b(X_t^\eta) + \eta F(X_t^\eta)) dt + \sqrt{\frac{2}{\beta}} dW_t,$$

where  $b, F : \mathbb{R}^d \rightarrow \mathbb{R}^d$  are smooth,  $F$  bounded,  $\beta > 0$ , and  $\eta \in \mathbb{R}$ . We assume that for each  $\eta$  it admits a unique invariant measure  $\nu_\eta$ .

## CONTRACTIVITY ASSUMPTION

We assume that there exists  $m > 0$  and  $M \geq 0$  such that

$$\langle x - y, b(x) - b(y) \rangle \leq -m |x - y|^2, \quad \text{if } |x - y| \geq M.$$

## TRANSPORT COEFFICIENTS

Estimator of linear response (observable  $R$  average 0 with respect to  $\nu_0$ )

$$\hat{\Phi}_{\eta,t} = \frac{1}{\eta t} \int_0^t R(X_s^\eta) ds \xrightarrow[t \rightarrow \infty]{\text{a.s.}} \alpha_{R,\eta} := \frac{1}{\eta} \int_{\mathbb{R}^d} R d\nu_\eta = \alpha_R + O(\eta)$$

**Asymptotic variance**  $\lim_{t \rightarrow \infty} t \text{Var}(\hat{\Phi}_{\eta,t}) = O(\eta^{-2})$

## COUPLING BASED ESTIMATOR

**Idea:** Use the reference dynamics to reduce the variance and bias of the estimator:

$$\hat{\Psi}_{\eta,t} = \frac{1}{\eta t} \int_0^t [R(X_s^\eta) - R(Y_s^0)] ds,$$

with  $(X_t^\eta, Y_t^\eta)_{t \geq 0}$  the solution of following system with coupled driving noises  $(W_t, \tilde{W}_t)_{t \geq 0}$

$$dX_t^\eta = (b(X_t^\eta) + \eta F(X_t^\eta)) dt + \sqrt{\frac{2}{\beta}} dW_t,$$

$$dY_t^0 = b(Y_t^0) dt + \sqrt{\frac{2}{\beta}} d\tilde{W}_t.$$

## SYNCHRONOUS COUPLING

Use the same Brownian motion to drive the two trajectories, i.e  $W = \tilde{W}$ . In this case the difference process is  $C^1$ :

$$d(X_t^\eta - Y_t^0) = (b(X_t^\eta) + \eta F(X_t^\eta) - b(Y_t^0)) dt.$$

If the drift  $b$  is contractive everywhere, i.e. there exists  $M = 0$ , then

$$|X_t^\eta - Y_t^0| \leq \left( |X_0^\eta - Y_0^0| - \frac{\eta \|F\|_\infty}{2m} \right) e^{-mt} + \frac{\eta \|F\|_\infty}{2m}.$$

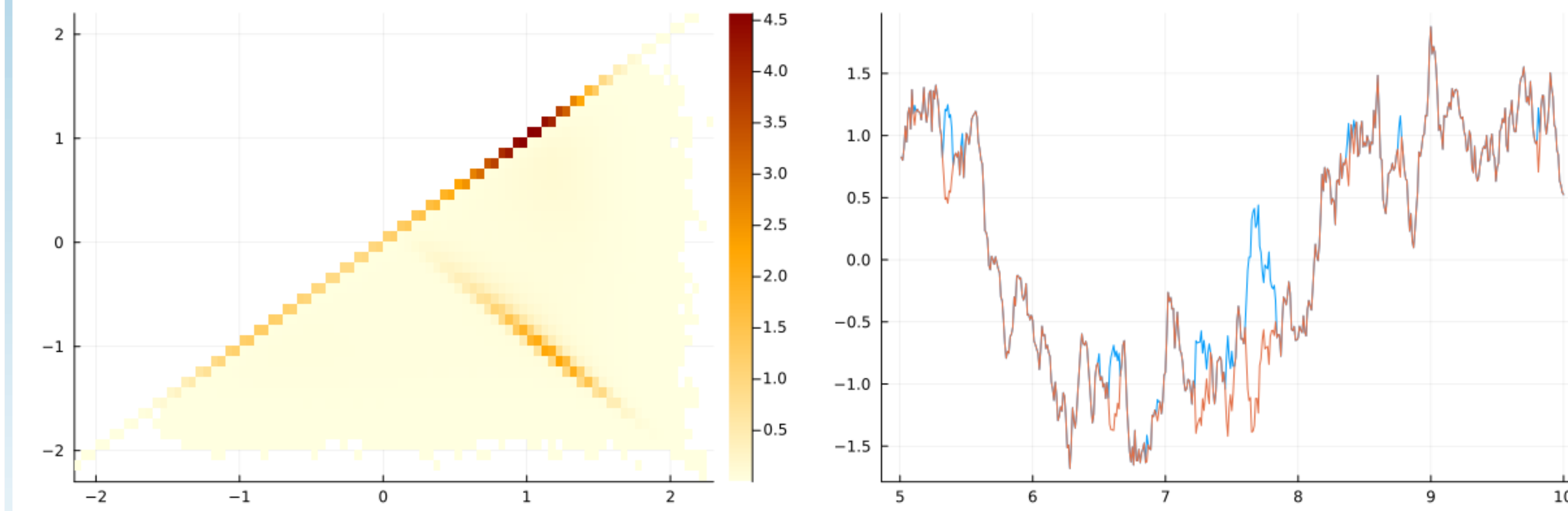
Thus, if  $X_0^\eta = Y_0^0$  then  $\mathbb{E} \left[ \left| \hat{\Psi}_{\eta,t}^{\text{sync}} \right|^p \right]$  is uniformly bounded as  $\eta \rightarrow 0$ .

## STICKY COUPLING

If the drift is *not* contractive everywhere, synchronous coupling can fail spectacularly because the drift does not necessarily bring together the trajectories.

**Idea:** Use the noise to bring the trajectories together.

Constructed in [4], uses reflection coupling to bring the trajectories together and is stick in the sense that  $|X_t^\eta - Y_t^0| \leq r_t^\eta$  a.s where  $(r_t^\eta)_{t \geq 0}$  is diffusion on  $[0, \infty)$  with a sticky boundary condition at zero



**Figure:** Sticky coupling of a 1D particle in a double well potential perturbed by a constant force to the right, i.e.  $b(x) = -4x(x^2 - 1)$  and  $\eta F(x) = 1$ . **Left:** histogram of coupled process; **Right:** segment of trajectory of coupled process

**Problem:** Sticky coupled process is highly degenerate. Zero set of  $r^\eta$  is a fat random Cantor set. Existence of invariant measure and its ergodic properties are unclear.

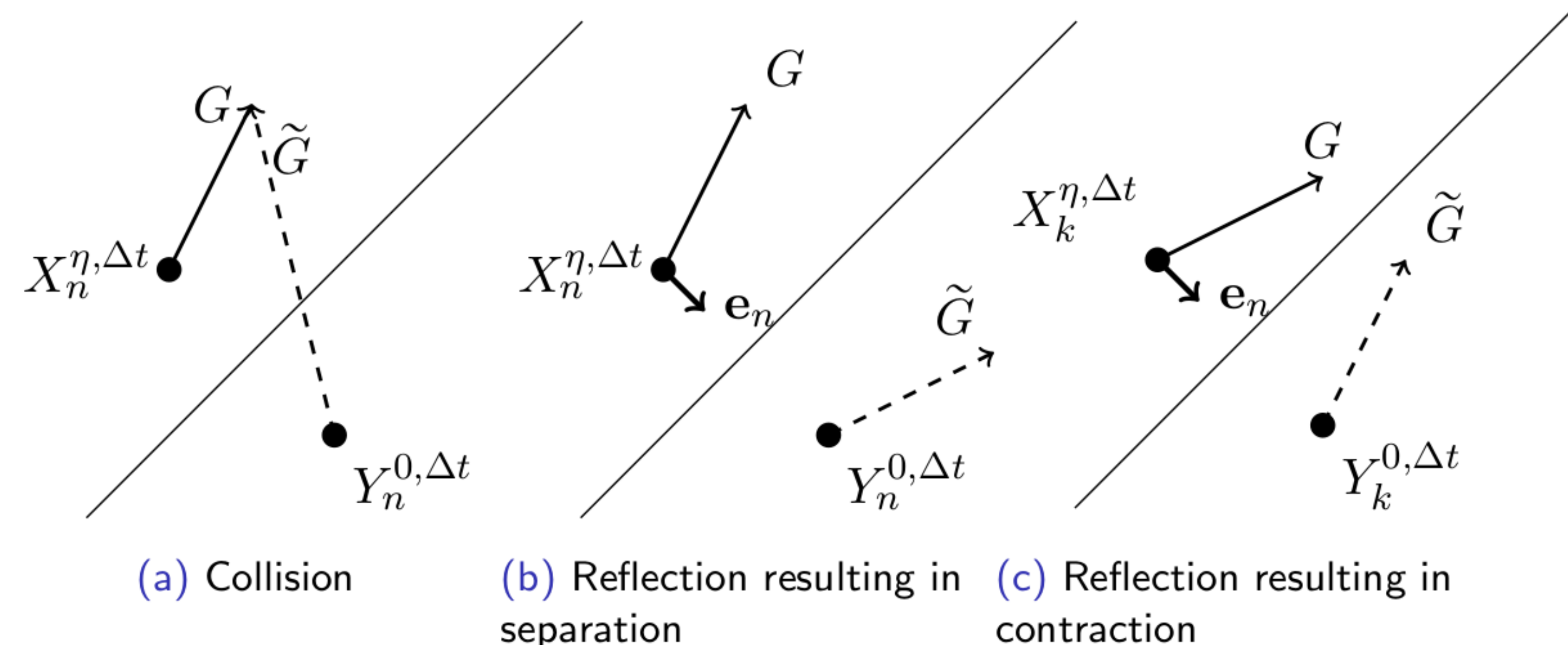
## DISCRETE STICKY COUPLING

**Work around:** Use discrete-time sticky coupling [2]:  $G_{n+1} \sim \mathcal{N}(0, \text{Id})$  and  $U_{n+1} \sim \text{Unif}([0, 1])$

$$X_{n+1}^{\eta, \Delta t} = X_n^{\eta, \Delta t} + \Delta t [b(X_n^{\eta, \Delta t}) + \eta F(X_n^{\eta, \Delta t})] + \sqrt{\frac{2\Delta t}{\beta}} G_{n+1},$$

$$Y_{n+1}^{0, \Delta t} = \begin{cases} X_{n+1}^{\eta, \Delta t}, & \text{if } U_{n+1} \leq p_{\Delta t, \beta}(X_n^{\eta, \Delta t}, Y_n^{0, \Delta t}, G_{n+1}), \\ Y_n^{0, \Delta t} + \Delta t b(Y_n^{0, \Delta t}) + \sqrt{\frac{2\Delta t}{\beta}} [\text{Id} - 2\mathbf{e}_n \mathbf{e}_n^T] G_{n+1} & \text{otherwise,} \end{cases}$$

with  $p_{\Delta t, \beta}(x, y, g)$  being the overlap of the two trajectories marginal transition densities at the  $n$ -th step.



The discrete-time sticky coupled process admits a unique invariant probability measure,  $\mu_{\eta, \Delta t}$  and is geometrically ergodic with respect to this measure. Furthermore, we can control the mass  $\mu_{\eta, \Delta t}$  puts off the diagonal:

$$\int_{\mathbb{R}^d \times \mathbb{R}^d} \mathbf{1}_{\{x \neq y\}} e^{-c(|x|^2 + |y|^2)} \mu_{\eta, \Delta t}(dx dy) \leq C\eta \left( \nu_{\eta, \Delta t}(e^{-c|x|^2}) + \nu_{0, \Delta t}(e^{-c|x|^2}) \right).$$

## DISCRETE-TIME ESTIMATOR

We introduce the following discrete-time estimator of  $\alpha_{R,\eta}$

$$\hat{\Psi}_{\eta, N}^{\text{sticky}, \Delta t} = \frac{1}{\eta N} \sum_{n=0}^{N-1} [R(X_n^{\eta, \Delta t}) - R(Y_n^{0, \Delta t})]$$

**Theorem 1** Let  $\eta_* > 0$  and  $R$  such that  $\nu_0(R) = 0$ . Assume that  $X^{\eta, \Delta t}$  and  $Y^{0, \Delta t}$  have the same initial value. Under the contractivity assumption and some technical assumptions, there exists  $K_1, K_2 > 0$  such that

$$\forall \eta \in [-\eta_*, \eta_*], \quad \lim_{N \rightarrow \infty} N \text{Var}(\hat{\Psi}_{\eta, N}^{\text{sticky}, \Delta t}) \leq \frac{K_1}{\eta},$$

and

$$\left| \mathbb{E}[\hat{\Psi}_{\eta, N}^{\Delta t}] - \alpha_{R, \eta} \right| \leq K_2 \left( \frac{1}{\Delta t N} + \Delta t \right).$$

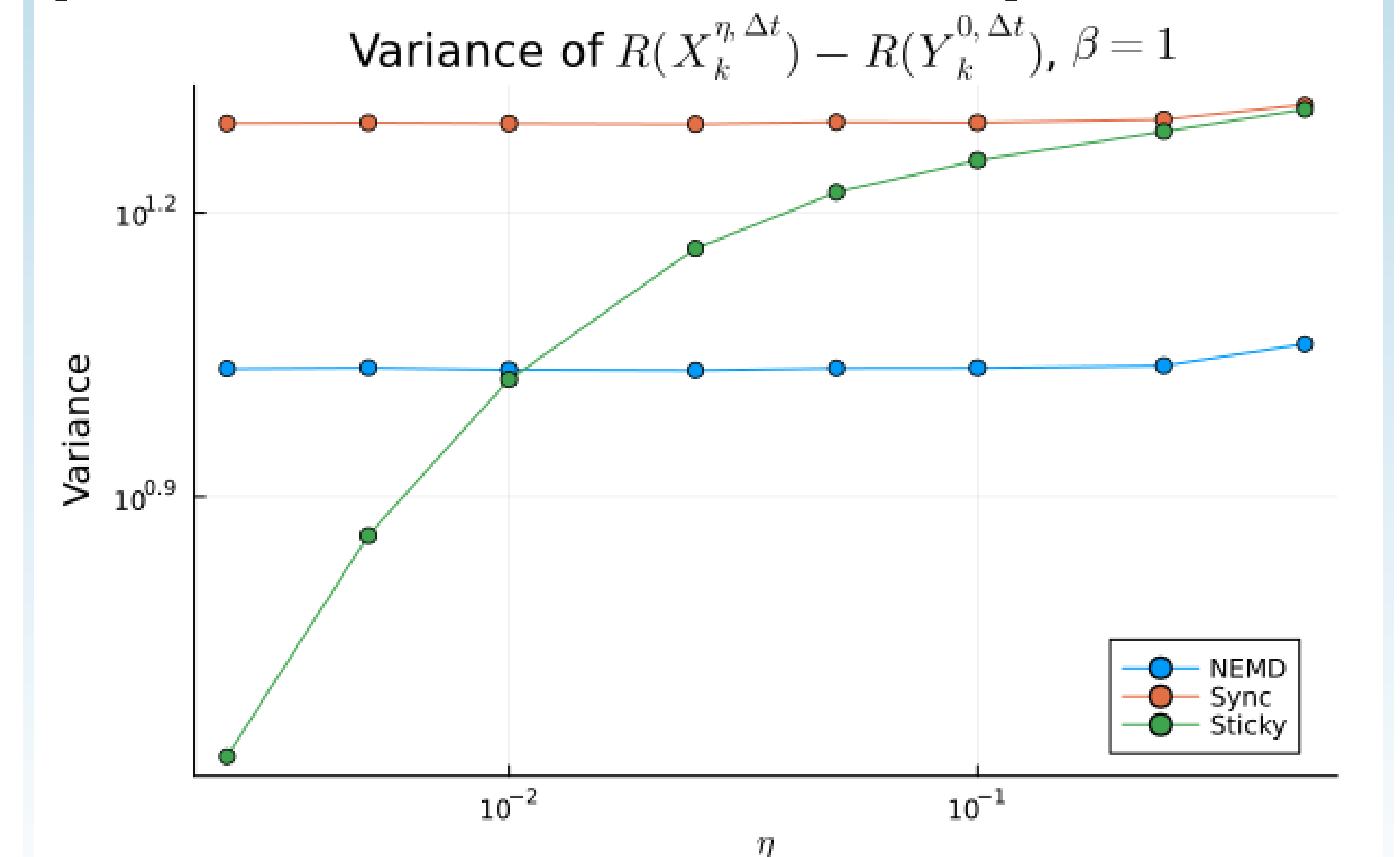
## NUMERICAL ILLUSTRATION

Example: Lennard-Jones Cluster in 2D:

The drift is given by  $b(x) = -\nabla(U_1 + U_2)$  with interaction part  $U_1$

$$U_1(x) = \sum_{i \geq j} \left[ \left( \frac{1}{|r_{ij}|} \right)^{12} - 2 \left( \frac{1}{|r_{ij}|} \right)^6 \right],$$

with  $r_{ij} = |x^i - x^j|$  if  $i < j$  and  $r_{ii} = |x^i|$ . And  $U_2$  a confining potential For  $F$  we use sine shear on each particle.



Code available: [github.com/shiva-darshan/sticky\\_coupling](https://github.com/shiva-darshan/sticky_coupling)

## REFERENCES

- [1] N. Bou-Rabee, A. Eberle and R. Zimmer Coupling and Convergence for Hamiltonian Monte Carlo 2020
- [2] A. Durmus, A. Eberle, A. Enfroy, A. Guillin, and P. Monmarché Discrete sticky couplings of functional autoregressive processes 2021
- [3] A. Eberle, A. Guillin, R. Zimmer Couplings and quantitative contraction rates for Langevin dynamics 2019
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- [5] Lelièvre, T. and Stoltz, G. Partial differential equations and stochastic methods in molecular dynamics. *Acta Numerica*, 2016.